

Study of Filament Features of Edge Plasma Fluctuations in Heliotron J

臧臨閣¹, 大島慎介², 西野信博³, 水内亨², 笠嶋慶純¹, 沙夢雨¹, 橋本 紘平¹, 向井清史⁴,
李炫庸¹, 長崎百伸², 岡田浩之², 南貴司², 小林進二², 山本聡², 史楠²,
木島滋², 中村祐司¹, 佐野史道²

ZANG Linge¹, S. OHSHIMA², N. NISHINO³, T. MIZUUCHI², K. KASAJIMA¹, M. SHA¹, K. HASHIMOTO¹,
et al.

京大エネ科¹, 京大エネ研², 広大院工³, 核融合研⁴
GSES Kyoto Univ.¹, IAE Kyoto Univ.², GSE Hiroshima Univ.³, NIFS⁴

To investigate plasma fluctuations and particle transport around LCFS region, a Langmuir probe array has been installed at #14.5 section in Heliotron J. In a condition-fixed ECH discharge series, supersonic molecular-beam (SMB) was injected at 222ms from #11.5 port. The Langmuir probe array was scanned along the radial direction shot by shot to get the radial profiles of ion saturation current (I_s) and floating potential (V_f).

The conditional average (CA) method is applied to \tilde{I}_s (\tilde{A} means 2kHz high pass filtered signal of A) data to extract the blob. The trigger condition of CA is set to 2.5σ in this case, where σ is the standard deviation of \tilde{I}_s during a time window of 5ms. The CA results of \tilde{I}_s , \tilde{E}_p and a blob-induced particle flux Γ_{blob} at two radial positions ($r-a = -4$ mm and $+10$ mm) are shown in Figs. 1(a) and (b). Here, The Γ_{blob} is calculated by $\tilde{I}_s \cdot \tilde{E}_p$, where E_p is the poloidal electric field evaluated from two floating potential signals of two electrodes arranged in the poloidal direction. The positive Γ_{blob} means the outward flux. The shapes of \tilde{I}_s , \tilde{E}_p and Γ_{blob} just

after SMBI (225-230 ms, red lines in Fig.1) do not change very much between $r-a = -4$ mm and $+10$ mm, here the blob count number per millisecond is about 3-4 for each position. On the other hand, long after SMBI (245-250 ms, blue lines), \tilde{E}_p -structure is different between two positions, $r-a = -4$ mm and 10mm. A characteristic structure, which is usually observed for the blob, is not observed at $r-a = -4$ mm, but, at $r-a = +10$ mm, the characteristic structure is observed. This and the data in-between the two position suggest the blobs are born probably near $r-a = +10$ mm in this timing (long after SMBI). The averaged particle flux carried by a single blob reduced to a small value compared to that at 225-230 ms, due to smaller values of \tilde{I}_s and \tilde{E}_p . The count number per millisecond, however, increased from 1 to 6 at this timing.

Figure 2 shows the radial profiles of time averaged I_s , E_r (radial electric field) and total Γ_{blob} during the time window (calculated from the CA results). In Fig. 2(a), the gradient of I_s at 245-250ms (long after SMBI) is much steeper than that at 225-230ms (just after SMBI). As in Fig. 2(c), however, the blob-induced particle flux at 245-250 ms is much reduced from that at 225-230 ms. The much steeper E_r profile in the region of -4 mm $< r-a < +6$ mm long after SMBI may be a candidate to explain the reduced blob-induced particle flux and increased count number. It can be expected that the sheared flow might tear up the coherent blob structures and reduced the blob-induced transport.

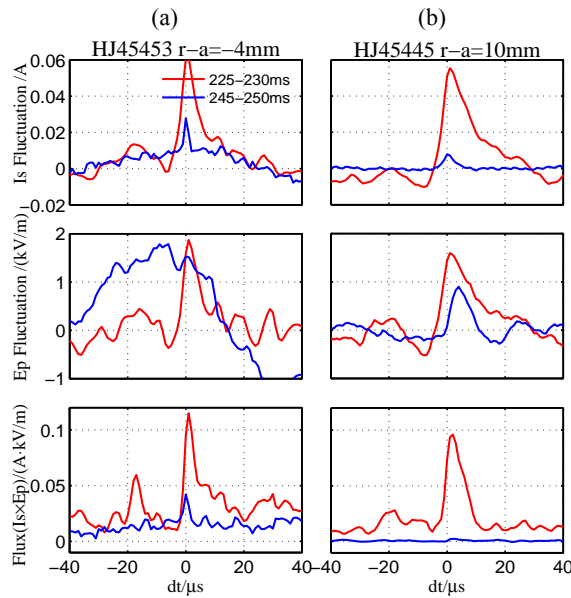


Figure 1. Auto CA of \tilde{I}_s , cross CA of \tilde{E}_p and blob induced particle flux; red: 225-230ms, blue: 245-250ms. Location: (a) $r-a=-4$ mm; (b) $r-a=10$ mm.

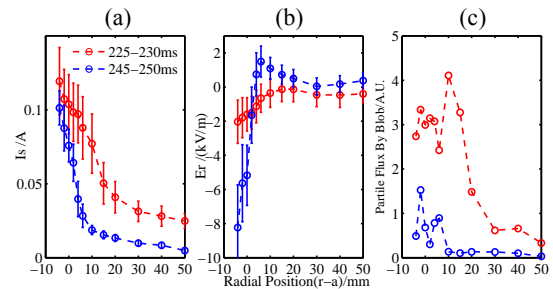


Figure 2. Radial profiles of (a) I_s ; (b) E_r ; (c) particle flux induced by blob. Red: 225-230ms, just after SMBI. Blue: 245-250ms.