

## 非平衡プラズマの温度に関する統計力学的検討

## Statistical-Mechanical Discussion on Temperatures of Non-Equilibrium Plasma

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It is well known that the electron energy distribution function (EEDF)  $F$  becomes far from a Maxwellian one in a weakly-ionized plasma [1]. However, chemical kinetics of counterparts of electron collisions are not often included in such calculations. We discussed the dissociation degree of oxygen molecule, with simultaneous solution of the two-term approximated Boltzmann equation [2]. Here, the definition of the “temperature” should be reconsidered from the statistical thermodynamic viewpoint, which has seldom been discussed. Let the internal energy  $U$  and the entropy  $S$  for electron gas in a weakly-ionized plasma. Then, the first law of thermodynamics leads:

$$\left(\frac{\partial S}{\partial U}\right)_{V, N_i} = \frac{1}{T}. \quad (1)$$

Here, we neglect change in its volume or electric potential. Hereafter,  $T$  in Eq. (1) is referred to as the thermodynamic electron temperature  $T_e^{\text{th}}$ . Meanwhile,

$$U = \langle \epsilon \rangle = \int_0^\infty \epsilon F(\epsilon) d\epsilon. \quad (2)$$

We can also calculate the Gibbs electron entropy as

$$S = -k \int_0^\infty F(\epsilon) \ln [F(\epsilon)] d\epsilon. \quad (3)$$

In the present study, we chose microwave discharge oxygen plasma in a cylindrical tube of 2.6 cm $\phi$  with its discharge pressure 1 Torr. We fixed the electron density  $N_e = 2 \times 10^{11} \text{cm}^{-3}$  and the gas temperature 1000 K.

Meanwhile, many researchers define the electron kinetic temperature  $T_e^{\text{k}}$  from  $U$  of Eq. (2) as follows [3]:

$$T_e^{\text{k}} = \frac{2}{3k} U, \quad (4)$$

since the kinetic temperature agrees with the equilibrium electron temperature for the plasmas with the Maxwellian EEDF. Although Alvarez *et al.* concluded that  $T_e^{\text{th}} = T_e^{\text{k}}$  [3], they never solved chemical kinetics of discharge gases

of the plasmas. We found the bulk ( $\epsilon \leq 8.4$  eV) and tail regions ( $\epsilon > 8.4$  eV), which has different slopes with their temperatures  $T_1$  and  $T_2$  ( $< T_1$ ), respectively. This is due to the electron-impact oxygen dissociation, which causes non-equilibrium for the EEPF.

Fig. shows a comparison between various temperatures defined to the oxygen plasma. It should be remarked that  $T_e^{\text{th}} \neq T_e^{\text{k}}$ , which does not agree with Alvarez *et al.*'s result [3]. It should be noted that the thermodynamic temperature  $T_e^{\text{th}}$  is more close to  $T_1$ , the electron temperature of the bulk component than  $T_e^{\text{k}}$  is. Further thermodynamic discussion is necessary to understand the physical meaning of each temperature, particularly for the relevance to Tsallis entropy.

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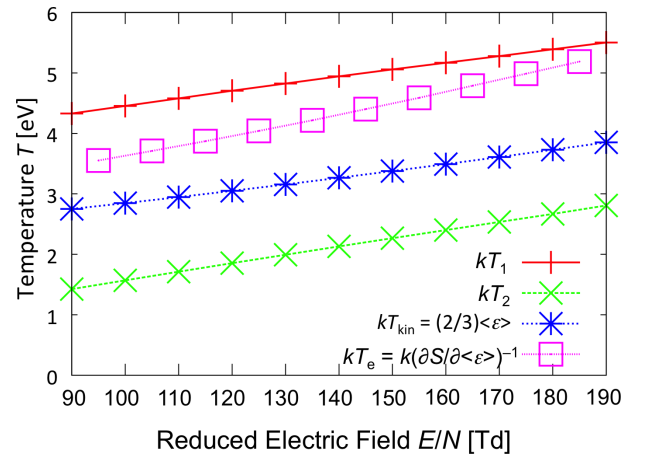


Fig. 1. Comparison between various electron temperatures defined in the present study for the oxygen plasma.

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[3] R. Alvarez, J. Cotrino and A. Palmero: EPL. **105** 15001 (2014).