原子構造計算プログラム FAC による多価タングステンイオンのスペクトル解析 Analysis of spectra from multiply ionized W ions by atomic structure code FAC

仲野友英¹、村上泉²、鈴木千尋² T. Nakano¹, I. Murakami², C. Suzuki²

> ¹原子力機構、²核融合研 ²JAEA, ²NIFS

Tungsten is one of the candidates for plasma facing components in future fusion devices such as ITER and DEMO. However, because of its high Z number, tungsten is not fully ionized in plasmas, and the intense line-radiation dissipates plasma energy, leading to a plasma collapse. Hence it is of significance to avoid accumulation of tungsten ions in the core plasma. For this purpose, it is required to establish a method to determine tungsten ion density in the plasma quantitatively. Spectroscopic measurement of tungsten spectra has been widely used and analysis with an atomic code has been extensively performed. In the present work, tungsten spectra measured in JT-60U and LHD have been analyzed with an atomic code, named FAC [1].

Figures 1 (a) and (b) show a vacuum ultra-violet spectra taken, respectively, from the peripheral plasma of JT-60U at $T_e \sim 2$ keV and from the core plasma of LHD at $T_e \sim 3$ keV. In order to identify tungsten lines, spectrum calculation with FAC was performed for W¹⁴⁺ - W⁷¹⁺. In the calculation, for instance, for W²⁷⁺ (Ag-like tungsten ion), the following electron configurations besides the ground state ($4d^{10}4f$) were considered:

 $4d^9 + 4f^2$

 $4d^{10}_{1}$ + 5s, 5p, 5d, 5f, 5g, 6s

 $4d^9 4f + 5s, 5p$

In calculation for $W^{14+} - W^{24+}$, only $\Delta n=0$ transitions between n=4 level were considered because calculation of $\Delta n=1$ transitions from n=5 to n=4 consumed huge computation time. Then, a collisional-radiative model, which included the processes of (de-) excitation between these levels and radiative transition from these levels, was used to calculate spectra of each W ion. As shown in Fig. 2, many spectral lines due to transitions between 4p - 4d and 4d - 4f, in particular from $W^{27+} - W^{34+}$, distributed between 4.5 nm and 5.5 nm. These spectra, weighted by a fractional abundance, were summed up and a synthesized spectrum was obtained.

In Fig 1 (a), comparison of the measured and the synthesized spectrum for JT-60U [2] is shown. The measured spectrum is well reproduced by the synthesized spectrum, which includes contribution of an additional peak of the fractional abundance around W^{32+} besides a peak around W^{43+} . This additional peak of the fractional abundance contributes to the broad peak observed around 5.0 nm in Fig. 1 (a). In contrast, because the synthesized spectrum for LHD does not include contribution of the



Fig. 1. Measured and synthesized spectrum (a) for JT-60U and (b) LHD.



additional peak, discrepancy between the measured and the synthesized spectrum is significant around 5.0 nm as shown in Fig. 1(b), showing necessity for adding the additional peak for lower charge states of W ions. This necessity is commonly seen in JT-60U and LHD. The reason why the additional peak is needed is unknown and will be investigated.

- [1] Gu, M. F. et al., Astrophys. J. 582 (2003) 1241.
- [2] Nakano, T. et al,: Nucl. Fusion. 49 (2009) 115024.