

## Analysis of Current Quench of the JT-60U tokamak by using two dimensional MHD equilibrium calculation Code (DINA).

2次元軸対称MHD平衡計算コードDINAを用いたJT-60Uにおける  
電流クエンチ解析

Syo Kawakami, Shibata Yoshihide, Watanabe Kiyomasa, Ohno Noriyasu,  
Kajita Shin, Okamoto Masaaki, Isayama Akihiko, Sugihara Masayoshi,  
Kawano Yasunori, V.E. Lukash, R.R. Khyarutdinov  
河上翔, 柴田欣秀, 渡邊清政<sup>1</sup>, 大野哲靖, 梶田信<sup>2</sup>, 岡本征晃<sup>3</sup>,  
諫山明彦<sup>4</sup>, 杉原正芳<sup>4</sup>, 河野康則<sup>4</sup>, V.E. Lukash<sup>5</sup>, R.R. Khayrutdinov<sup>5</sup>

Graduate School of Engineering, Nagoya University, Furou-cho, Chikusa-ku, Nagoya, Aichi 464-8603, Japan  
名古屋大学工学研究科, 〒464-8603 名古屋市千種区不老町

<sup>1</sup>National Institute for Fusion Science, Oroshi-cho, Toki, Gifu 509-5292, Japan  
自然科学研究機構 核融合科学研究所, 〒509-5292 岐阜県土岐市下石朝 322-6

<sup>2</sup>Eco Topia Science Institute, Nagoya University, Furou-cho, Chikusa-ku, Nagoya, Aichi 464-8603, Japan  
名古屋大学エコトピア研究機構, 〒464-8603 名古屋市千種区不老町

<sup>3</sup>Ishikawa National College Technology, Kitachoujo, Tsubata-cho, Kahoku-gun, Ishikawa 929-0392, Japan  
石川工業高等専門学校, 〒929-0392 石川県河北郡津幡町北中条夕 1

<sup>4</sup>Japan Atomic Energy Agency, Muramatsu, Tokai-mura, Naka-gun, Ibaraki 319-1184, Japan  
日本原子力研究開発機構, 〒319-1184 茨城県那珂郡東海村村松 4-49

<sup>5</sup>Kurchatov Institute, Academician Kurchatov Sq 1, Moscow, Russia 123 098

During the disruption of tokamak, the plasma current rapidly decays to zero and it gives the large mechanical force in the structure. In this study, we predicted the time evolution of plasma current during the disruption by using the disruption simulation code, which is called DINA, and investigated the effects of electron temperature,  $T_e$  and an effective charge,  $Z_{\text{eff}}$  on the predicted time evolution of plasma current,  $I_p$  in JT-60U radiative disruption. In the calculation, it was found that time evolution of  $I_p$  is good agreement with the experimental behavior of  $I_p$  when we assumed  $T_e=30\text{eV}$  and  $Z_{\text{eff}}=3$ . However, internal inductance, which indicates the shape of current density distribution, evaluated by DINA code differs from the experimental values. Hence, we need to match the experimental and calculated results in order to predict correctly the plasma behavior during the disruption.

### 1. Introduction

The plasma current and cross-section plasma rapidly decay to zero during tokamak disruption. They generate a halo current and an eddy current, and such currents give a large mechanical force into the structure around the plasma. It is a significant problem for stable operation of fusion tokamak device such as ITER.

In this study, we simulate the behavior of the plasma current during the radiative disruptive discharge in JT-60U by using a disruptive simulation code, DINA in order to clarify the determination mechanism of plasma current decay during the current quench [1].

### 2. Modification of DINA

DINA is two dimensional MHD equilibrium calculation code for disruption simulation. The rate of plasma current quench in DINA, taking the time evolution of the poloidal flux into account, is determined with specified values of electron temperature,  $T_e$ , and an effective charge,  $Z_{\text{eff}}$  [2,3].

$T_e$  and  $Z_{\text{eff}}$  are calculated by a Power Balance equation. However, the upper limit of  $T_e$  calculation result sets to 10eV because  $T_e$  has the single unique solution. To investigate the effect of those plasma parameters on the time evolution of plasma current during disruption, we analyze in the radiative disruptive disruption of JT-60U. Table I shows the assumption values of  $T_e$  and  $Z_{\text{eff}}$  in this calculation [4]. In case A, the  $T_e$  and the  $Z_{\text{eff}}$  were calculated by DINA code. In case B, we assumed  $T_e$  equals 20eV and  $Z_{\text{eff}}$  calculated from Power Balance equation. In case C, we assumed  $Z_{\text{eff}}$  equals 3.

Table I. The assumption values of  $T_e$  and  $Z_{\text{eff}}$  in this calculation

	$T_e$ (eV)	$Z_{\text{eff}}$
Case A	—	—
Case B	20(constant)	—
Case C	20(constant)	3
Case D	30(constant)	3

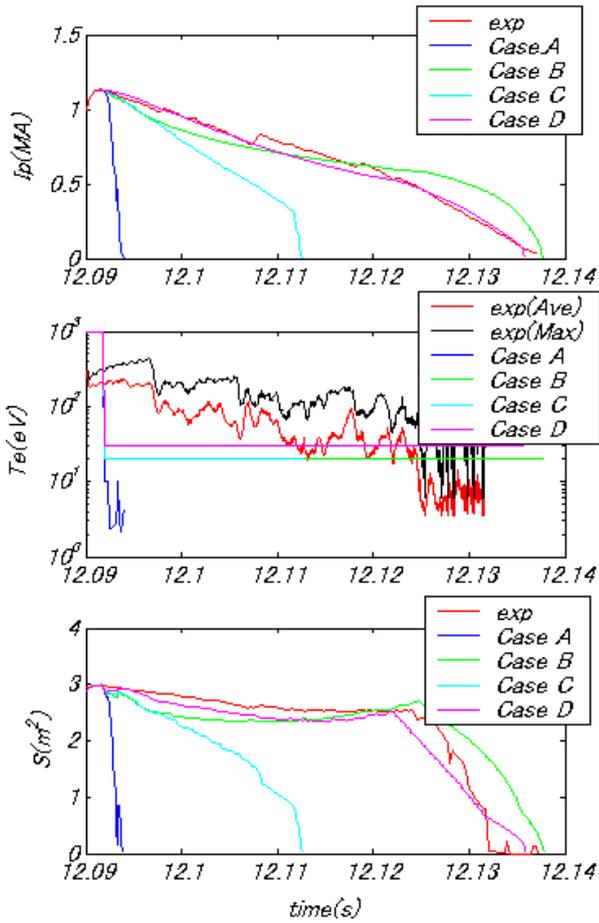


Fig.1. Time evolutions of plasma current  $I_p$ , electron temperature  $T_e$ , and plasma cross-section  $S$

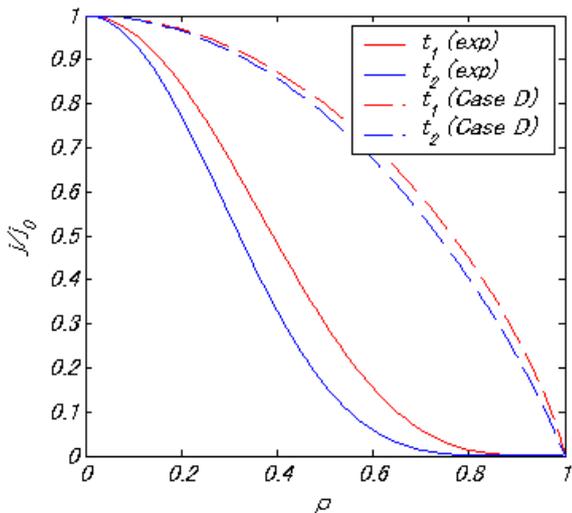


Fig.2. Profiles of normalized current density distribution

### 3. Result

Waveforms of time evolutions of plasma current  $I_p$ ,  $T_e$ , and plasma cross-section,  $S$ , are shown in Fig.1. In case A, DINA does not reproduce behavior of the plasma current, because  $T_e$  decays rapidly compared with the experimental data. In case D, time evolutions of the  $I_p$  and  $S$  were good agreement with the experimental data.

Fig.2 shows profiles of normalized current density distributions at  $I_p = I_{p0}(t_1)$  and  $I_p = 0.9 \times I_{p0}(t_2)$ . Here  $I_{p0}$  is the value of plasma current just after current quench starts.  $\rho$  is a normalized plasma small radius. The current density profile is calculated by the Eq(1)

$$j/j_0 = (1 - \rho^2)^v, \quad (1)$$

$$v = (e^{l_i} - 1.65)/0.89, \quad (2)$$

where  $v$  is the peaking factor and  $l_i$  is the internal inductance [3]. In experimental data, we calculated  $v$  using the Eq(2). When  $I_p = I_{p0}$ , in case D, this profile does not match the experimental measurement. Furthermore, time change of experimental current density profile is peaking at center relative to time change of DINA.

In addition to, in case C,  $I_p$  decays more rapidly relative to case B and in case B, averaged  $Z_{eff}$  is 1.29. Therefore, the rate of plasma current quench depends on  $Z_{eff}$ .

### 4. Summary

We analyzed the current quench in the radiative disruption of JT-60U by using DINA. When we assumed  $T_e = 30\text{eV}$  and  $Z_{eff} = 3$  in DINA, we showed the behavior of the plasma current is good agreement and  $Z_{eff}$  affects to  $I_p$ . However, DINA profile of normalized current density distribution at  $I_p = I_{p0}$  does not match the experimental measurement, because flat  $T_e$  profile was assumed in the present DINA calculation.

For future work, we will introduce the measured  $T_e$  profile to the DINA code to reproduce the profile of the normalized current density distribution more accurately.

### References

- [1] R.R.Khayrutdinov and V.E.Lukash, J. Comput. Physics, **109**(1993) 193-201.
- [2] H. Ohwaki, *et al.*, Plasma and Fusion Research Regular Articles, **1**(2006) 016.
- [3] J.A.Wesson, Tokamaks 3rd edition, Oxford University Press, (2004) Ch 3,6,7.
- [4] Y.Shibata, *et al.*, Nucl Fusion, **50**(2010)025015.