

Generation of Quasi-Pair Fullerene Ion Plasma for Purification of Lithium Endohedral C₆₀

リチウム内包C₆₀高純度化のための準ペアラーレンイオンプラズマ生成

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The plasma consisting of the lithium endohedral fullerene positive ion (Li@C₆₀⁺) and the fullerene negative ion (C₆₀⁻) is generated by an electron beam ionization method for purification of Li@C₆₀. The ionization potential of Li@C₆₀ is measured for the first time and it is found to be lower than that of C₆₀. In addition, it is implied that argon positive ions (Ar⁺) can dissociate the Li@C₆₀ - C₆₀ cluster more effectively than the high energy electron beam, resulting in generation of high density Li@C₆₀⁺ - C₆₀⁻ plasma.

1. Introduction

Alkali metal endohedral fullerenes, such as lithium endohedral fullerene (Li@C₆₀), have attracted much attention [1] because they are expected to be used in the field of single molecule electronics, owing to their unique electronic properties which have roots from the lithium off-center position in the C₆₀ cage [2].

Although Li@C₆₀ has been isolated from C₆₀ by a chemical method [3], this method takes a long time and needs to chemically modify the surface of Li@C₆₀, which makes it difficult to be used for a practical application. In this sense, a plasma separation method is expected as one effective way to purify the Li@C₆₀ sample which is mass-produced by Ideal Star Ink. in collaboration with Tohoku University, because the plasma process is fast and simple, and no chemical modification is needed during the whole process. However, the problem of this method is that both the C₆₀ and Li@C₆₀ can become positive ions [4], though an ideal Li@C₆₀⁺ - C₆₀⁻ plasma is required for the plasma separation. Therefore, the difference in the plasma characteristics, e.g., electron impact ionization cross section, between Li@C₆₀ and C₆₀ [5] should be found out to selectively generate Li@C₆₀⁺.

2. Experimental Setup

The fullerene plasma is generated by an electron beam ionization. The plasma source setup is schematically shown in Fig. 1. Thermionic electrons are emitted from a tungsten cathode by applying a DC power P_k and are accelerated by a potential difference between the cathode V_k and a grounded grid located at z = 0 mm. The mixture of

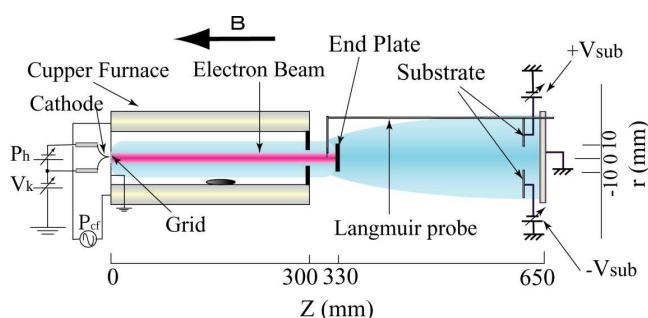


Fig. 1: The schematic illustration of experimental set up.

C₆₀ and Li@C₆₀ in a copper furnace is sublimated by heating up to 650 °C. High energy electrons easily ionize Li@C₆₀ molecules into Li@C₆₀⁺ while low energy electrons produced as a result of positively ionizing Li@C₆₀ are attached on C₆₀ molecules. The Li@C₆₀⁺ - C₆₀⁻ plasma is generated radially outside of the electron beam by a magnetic filter effect. An end plate is located at z = 330 mm, which cuts the electron beam and thus produces a pure Li@C₆₀⁺ - C₆₀⁻ plasma behind it. Stainless steel substrates are located at z = 650 mm. The C₆₀ and Li@C₆₀ ions are deposited on the substrates by applying positive and negative bias voltages to the substrates (V_{sub}). The deposited materials are measured by laser-desorption time-of-flight mass spectrometer (LD TOF-MS).

3. Experimental Results

Figure 2 shows the typical Langmuir probe voltage V_p - probe current I_p characteristics of Li@C₆₀⁺ - C₆₀⁻ plasma measured at z = 390 mm. Remarkably, positive and negative saturation currents are almost the same unlike the ordinary electron - ion plasma and both the floating potential

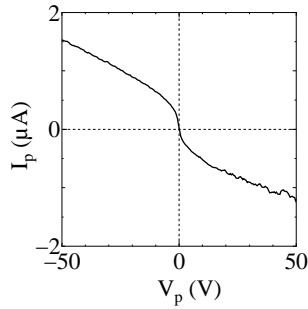


Fig. 2: The typical $V_p - I_p$ characteristics of $\text{Li@C}_{60}^+ - \text{C}_{60}^-$ plasma

and space potential are almost zero because the positive and negative ions weights are the same.

Figure 3 presents positive probe currents (I_p^+) of (a) C_{60} plasma and (b) $\text{Li@C}_{60}^+ - \text{C}_{60}^-$ plasma as a function of electron energy (ϵ_e) using the probe located at $z = 370$ mm. Here, the cathode current is constant (the current between the cathode and the grid: $I_k = 1$ mA) and probe bias voltage for I_p^+ and negative probe current (I_p^-) are -100 V and $+100$ V, respectively. It is noteworthy that there is a difference of threshold energy (ϵ_{th}) for second peak between C_{60} and $\text{Li@C}_{60}^+ - \text{C}_{60}^-$ plasmas. The ionization potential of the Li@C_{60} is approximately 3 eV lower than that of C_{60} . This result might be very useful in terms of the selective generation of Li@C_{60}^+ and C_{60}^+ .

Figure 4 gives the change of I_p^+ and I_p^- as a function of ϵ_e (a) without and (b) with introduction of the Ar plasma in the furnace area. It is shown that the plasma density in the diffusion area is increased significantly by the Ar plasma introduction, which is considered to dissociate the cluster of $\text{Li@C}_{60} - \text{C}_{60}$ inside the furnace. The $\text{Li@C}_{60} - \text{C}_{60}$ cluster consists of one Li@C_{60} molecule and 10 - 12 C_{60} molecules. The cluster can not keep its ion state for a long time because its Larmor radius is much larger than the device as explained in Fig.1. However, when the cluster is dissociated, one cluster is divided into more than 10 molecules and each molecule becomes the ionic state. The formed ions can diffuse behind the end

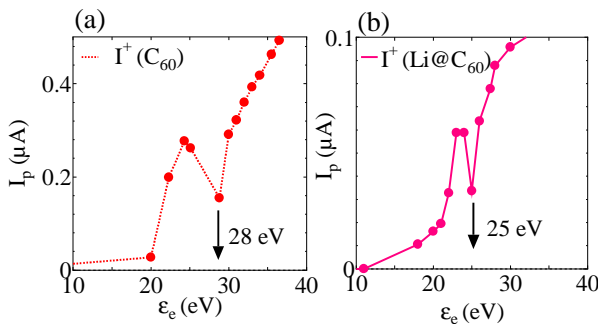


Fig. 3: The $\epsilon_e - I_p$ characteristics of (a) C_{60} plasma (b) $\text{Li@C}_{60}^+ - \text{C}_{60}^-$ plasma

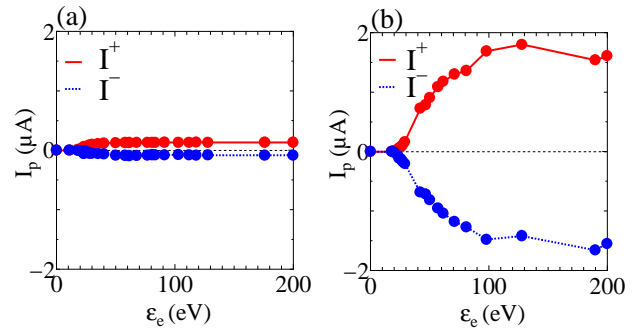


Fig. 4: $\epsilon_e - I_p$ characteristics of Li@C_{60} compound (a) with and (b) without Ar plasma.

plate. The content described above is the mechanism that the Ar plasma can increase the $\text{Li@C}_{60}^+ - \text{C}_{60}^-$ plasma density.

4. Summary

The $\text{Li@C}_{60}^+ - \text{C}_{60}^-$ plasma is successfully generated, which is confirmed by the Langmuir probe. The ionization potential of Li@C_{60} is found to be lower than that of C_{60} . It is the really valuable result in terms of realizing the selective ionization of Li@C_{60} from C_{60} and more efficient purification of Li@C_{60} . Furthermore, the Ar plasma is found to be useful to increase the density of Li@C_{60} plasma because the Ar plasma dissociates the $\text{Li@C}_{60} - \text{C}_{60}$ cluster more effectively than the electron beam. Our current result is very important for realizing the separation of Li@C_{60} and C_{60} . More detailed experiments are required to understand the ionization mechanism of Li@C_{60} . The deposited material on the substrate will also be studied as a future work.

References

- [1] K. Ohno, Y. Maruyama, K. Esfarjani, Y. Kawazoe, N. Sato, R. Hatakeyama, T. Hirata, and Michio Niwano: Phys. Rev. Lett. **76** (1996) 3590.
- [2] M. Pavanello, A. F. Jalbout, B. Trzaskowski, and L. Adamowicz: Chem. Phys. Lett. **442** (2007) 339.
- [3] S. Aoyagi, E. Nishibori, H. Sawa, K. Sugimoto, M. Takata, Y. Miyata, R. Kitaura, H. Shinohara, H. Okada, T. Sakai, Y. Ono, K. Kawachi, K. Yokoo, S. Ono, K. Omote, Y. Kasama, S. Ishikawa, T. Komuro, and H. Tobita: Nature Chem. **2** (2010) 678.
- [4] W. Oohara and R. Hatakeyama: Phys. Rev. Lett. **91** (2003) 205005.
- [5] S. Matt, B. Dunser, M. Lezius, H. Deutsch, K. Becker, A. Stamatovic, P. Scheier, and T. D. Mark: J. Chem. Phys. **105** (1996) 1880.