

Micro- and Macro-Scale Self-Organization in a Dissipative Plasma

ŠKORIĆ Miloš M.^{1,2}, SATO Tetsuya¹, MALUCKOV Aleksandra³ and JOVANOVIĆ Moma S.⁴

¹Theory and Computer Simulation Center, National Institute for Fusion Science, Toki-shi, 509-5292, Japan

²Vinča Institute of Nuclear Sciences, P.O.B. 522, 11001 Belgrade, Yugoslavia

³The Graduate University for Advanced Studies, Hayama, Kanagawa 240-01, Japan

⁴Department of Physics, University of Niš, P.O.B. 91, 18001 Niš, Yugoslavia

(Received: 11 December 1998 / Accepted: 5 August 1999)

Abstract

We study a nonlinear three-wave interaction in an open dissipative model of stimulated Raman backscattering in a plasma. A hybrid kinetic-fluid scheme is proposed to include anomalous kinetic dissipation due to electron trapping and plasma wave breaking. We simulate a finite plasma with open boundaries and vary a transport parameter to examine a route to spatio-temporal complexity. An interplay between self-organization at micro (kinetic) and macro (wave/fluid) scales is revealed through quasi-periodic and intermittent evolution of dynamical variables, dissipative structures and related entropy rates. At this point, a consistency with a general scenario of self-organization is found.

Keywords:

complexity, self-organization, stimulated Raman, scattering

1. Introduction

Self-organization (SO) is a generic process which describes a spontaneous formation of an ordered structure in a nonlinear far-from equilibrium system. Energy pumping, nonlinear instability, entropy production and expulsion are key governing processes. In this paper, an attempt is made to study kinetic self-organization [1] through a process of stimulated Raman backscattering (SRBS) in an underdense plasma [2]. Stimulated Raman scattering is a resonant three-wave (3WI) parametric instability which corresponds to the decay of an incident electromagnetic pump wave into a scattered wave plus an electron plasma wave (epw). Important effects, especially in the context of laser fusion schemes [3], are input energy loss, plasma heating and generation of suprathermal (hot) electrons. To emulate these effects we apply a hybrid three-wave phenomenological kinetic model of SRBS, proposed by

Škorić *et al.* [4].

2. A Hybrid Model

The basic equations are coupled three-wave equations for slowly varying complex amplitudes of a pump (a_0), backscattered wave (a_1) and epw (a_2), respectively:

$$\frac{\partial a_0}{\partial t} + V_0 \frac{\partial a_0}{\partial x} = -a_1 a_2,$$

$$\frac{\partial a_1}{\partial t} - V_1 \frac{\partial a_1}{\partial x} = -a_0 a_2^*,$$

$$\frac{\partial a_2}{\partial t} + V_2 \frac{\partial a_2}{\partial x} + \Gamma a_2 + i\sigma |a_2|^2 a_2 = \beta_0^2 a_0 a_1^*. \quad (1)$$

Parameter Γ is composed of collisional damping (Γ_{coll}), 'linear' (Landau) damping by hot – resonant electrons

Corresponding author's e-mail: skoric@nifs.ac.jp

($\gamma_1 \sim n_h$), where $n_h(t)$ stands for spatially averaged hot electron density:

$$\frac{dn_h(t)}{dt} = \frac{n_b(L, t)}{L} \int_{v_h - v_{tr}(L, t)}^{v_h + v_{tr}(L, t)} f_b dv - a \frac{v_h(t) n_h(t)}{L}, \quad (2)$$

(n_b is bulk electron density, f_b is bulk distribution function, v_h is velocity of hot electrons which equals the *epw* phase velocity, v_{tr} is trapping velocity and a is particle transport parameter) and nonlinear damping due to trapped resonant bulk electrons (in the thermal Maxwellian):

$$2\gamma_{nl} W(t) = \frac{mn_b(L, t)}{2L} \int_{v_h - v_{tr}(L, t)}^{v_h + v_{tr}(L, t)} v^3 f_b(v) dv, \quad (3)$$

where $W(t)$ is spatially averaged *epw* energy.

Open boundaries and re-emittance of fresh ambient electrons [5] is assumed. The particles (bulk and hot electrons) and energy (wave and particle) are being exchanged between a plasma and an environment through open boundaries with a conservation of particle number and total energy in the system. Accordingly, heat balance equation (effect of plasma heating; E_b , E_h and $\Phi_{tot} \rightarrow \Phi_b + \Phi_h + \Phi_q$ are the average bulk and hot electron energy or corresponding energy flux of the bulk, hot and return ambient electrons, respectively) takes form:

$$\frac{dW(t)}{dt} = \frac{d(E_b(t) + E_h(t))}{dt} + k \cdot \Phi_{tot}|_0. \quad (4)$$

3. Dissipative Structures and Entropy Rate

The simulation is performed via the central-difference numerical code [2], where the simulation parameters are: $n_0 = 0.1 n_{cr}$, $T_{b0} = 0.5$ keV and pump intensity, $\beta_0 = 0.0253$. Openness of a system, k (0-1) was chosen as a bifurcation parameter. Self-organization in strongly nonlinear far-from-equilibrium systems leads to a creation of ordered states that reflect an interaction of a given system with its environment. These novel dynamical structures or patterns, named dissipative structures to stress the crucial role of dissipation in their creation, have become a central theme of the science of complexity [1,2,4]. On the other hand, there is a fundamental role of the entropy, in particular, the rate of entropy change in an open system. The rate of entropy production and its removal basically governs self-organization features of a system.

First, we focus at self-organized dissipative structures developed at the macro-scale. Indeed, in our model, basic wave and fluid density variables were assumed to vary slowly in space-time. Therefore, we

expect that original spatio-temporal profiles, found in simulations, should correspond to large dissipative structures, self-organized at macro-scale levels. As an illustration, we plot the plasma wave profiles (Fig. 1), in particular, to reveal a genuine spatio-temporal nature of an intermittent regime as compared to regular dynamical regimes of the steady-state and quasi-periodic type [3,4]. Spatio-temporal complexity of quasi-steady and travelling wave patterns with regular and chaotic features is found in different states of self-organization.

Further, in Fig. 2 we plot the entropy rate $dS(t)/dt$ in time together with a spatio-temporal profile of the scattered wave energy. For an intermittent regime,

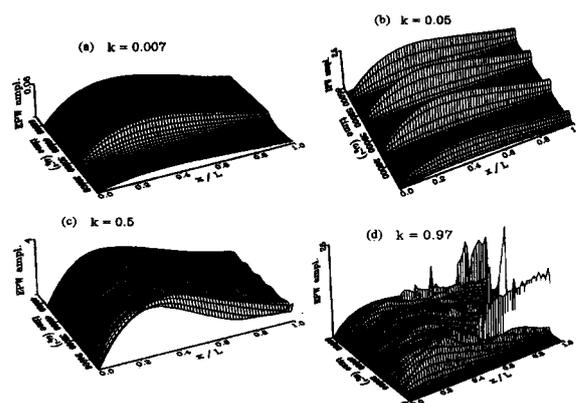


Fig. 1 Spatio-temporal profiles of the electron plasma wave for varying transport parameter k values. Different dissipative structures are seen on the route to complexity, from the steady-state via quasi-periodic to intermittent regimes.

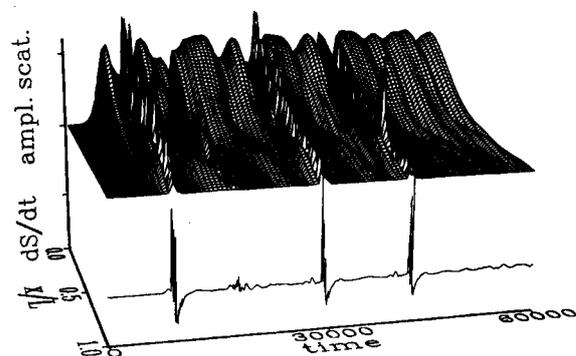


Fig. 2 Intermittent dissipative backscattered wave structures versus the corresponding entropy rate in time. Positive entropy jump coincides with an onset of chaos, while a negative burst indicates a transition from a chaotic to a laminar phase of SO at macro-scales.

featuring an interchange between chaotic and laminar phases, we find a clear evidence of structural transitions corresponding to the maximum (positive) and minimum (negative) entropy rate. As a striking example of self-organization in an open system we find a rapid entropy jump which coincides with an onset of a chaotic phase. Subsequent anomalous dissipation and entropy growth is halted by a sudden entropy expulsion into the environment. Negative burst in entropy rate indicates a bifurcation from a chaotic, back to a laminar quasi-periodic phase. Intervals of near zero entropy rate during a laminar phase, mean a net balance between the entropy production and its expulsion. This appears to be an example of a stationary nonequilibrium state possibly realized in a strongly nonlinear open system [1].

A hybrid nature of our model allows us to recover kinetic properties of self-organization. In Fig. 3 we see a three-dimensional view of the electron velocity distribution in time for different saturated Raman regimes, as indicated by values of parameter k . Kinetic self-organization of varying complexity is revealed in thermal and suprathermal (hot) regions of the electron distribution.

4. Summary

In summary, we believe that our findings appear to be a first indication of a generic intermittent scenario in a kinetic self-organization of anomalous Raman instability. At this point we may note that one is able to claim a consistency with the working hypothesis and general scenario of self-organization in plasmas [1,4]. As a further step, we expect an important justification of our hybrid-modeling of saturated Raman complexity by the novel open boundary particle simulation code, currently under development [5]. As an early illustration, we show in Fig. 4b, recent particle-in-cell simulation data for a model of an isolated plasma slab in a vacuum [6]. For same plasma parameters, particle simulations (Fig. 4b) show an evident support of above Raman reflectivity pattern, obtained for a closed ($k = 0.007$) system (Fig. 4a).

Acknowledgements

We acknowledge the partial support and the visiting professorship granted to one of us (M.M.Š.) by the Ministry of Education, Science and Culture of Japan. The hospitality of the National Institute for Fusion Science is gratefully acknowledged. This work was supported in parts by the Ministry of Science and Technology of the Republic of Serbia, Contract 01E11.

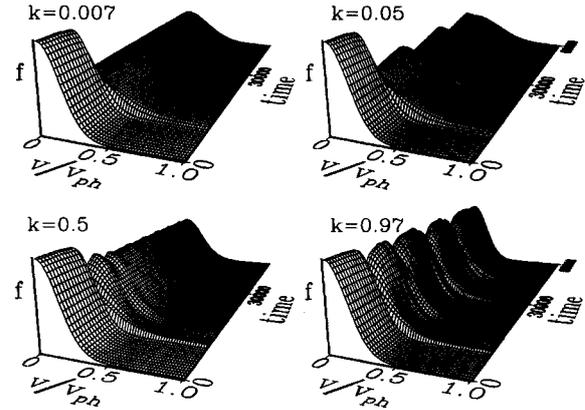


Fig. 3 Three-dimensional view of the electron velocity distribution in time for different saturated Raman regimes, as indicated by values of parameter k . Micro-kinetic scale self-organization of varying complexity is revealed in both thermal and suprathermal (hot) regions of the electron distribution.

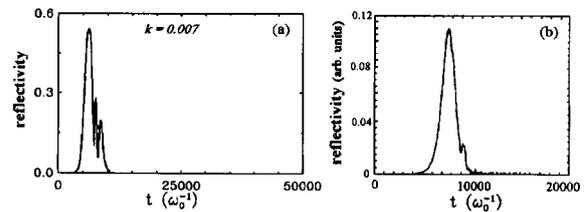


Fig. 4 Raman reflectivity versus time $[\omega_0^{-1}]$ for $k = 0.007$ from the hybrid model simulation (a) and the corresponding data (b) obtained by a $1\frac{1}{2}$ relativistic particle-in-cell code (after Miyamoto *et al.* [6]). The pump strength was $\beta_0 = 0.0253$ and the plasma parameters: $n_0 = 0.1 n_{crit}$, $T_e = 0.5$ keV, $L = 100 c/\omega_0$.

References

- [1] T. Sato and the Complexity Simulation Group, *Phys. Plasmas* **3**, 261 (1996).
- [2] M.M. Škorić, M.S. Jovanović and M. Rajković, *Phys. Rev. E* **53**, 4056 (1996).
- [3] W.L. Kruer, *The Physics of Laser Plasma Interactions*, (Addison, N.Y., 1988).
- [4] M.M. Škorić, T. Sato, A. Maluckov and M.S. Jovanović, NIFS Report, No.549 (1998); *Phys. Rev. E* (to appear).
- [5] H. Takamaru, T. Sato, R. Horiuchi, K. Watanabe and Complexity Simulation Group, *J. Phys. Soc. Jpn.* **66**, 3826 (1997).
- [6] S. Miyamoto, K. Mima, M. Škorić and M.S. Jovanović, *J. Phys. Soc. Jpn.* **67**, 1281 (1998).